

Approximation rates matter: sharp estimates for Target Measure Diffusion Maps and applications to molecular dynamics

Target Measure Diffusion Maps

Goal: Approximate infinitesimal generator \mathcal{L} for the time-reversible dynamics governed by the SDE

(1)
$$dX_t = \beta^{-1} \nabla \log(\mu)(x) dt + \sqrt{2\beta^{-1}} dW$$

We recall that the generator is given by

(2)
$$\mathcal{L} = \beta^{-1} \left(\Delta - \nabla \log(\mu) \cdot \nabla \right).$$

Input: Dataset $\mathcal{X} \subset \mathcal{M} \sim d\rho$, target measure μ , kernel bandwidth ϵ , and kernel matrix

(3)
$$[K_{\epsilon}]_{ij} = \exp(\epsilon^{-1} ||x_i - x_j||_2^2) := k_{\epsilon}(x_i, x_j)$$

Step 1 (Kernel Density estimation):

(4)
$$[M]_{ij} = \delta_{ij}\mu(\mathbf{x}_i)$$
 $D_{\epsilon} = n^{-1}\text{diag}(K_{\epsilon}\mathbf{1})$

Step 2 (Target measure renormalization):

(5)
$$K_{\epsilon,\mu} = K_{\epsilon} D_{\epsilon}^{-1} M^{1/2}$$

Step 3/Output: Generator:

$$D_{\epsilon,\mu} = n^{-1} \operatorname{diag}(K_{\epsilon,\mu}\mathbf{1}), P_{\epsilon,\mu} = D_{\epsilon,\mu}^{-1}K_{\epsilon,\mu}$$
(6)
$$L_{\epsilon,\mu}^{(n)} = \frac{I - P_{\epsilon,\mu}}{\epsilon}$$

Banisch et al. [1] prove that $4L_{\epsilon,\mu}$ is a Monte **Carlo approximation of** \mathcal{L} **:**

(7)
$$4\beta^{-1}L_{\epsilon,\mu}^{(n)}f(x) \rightarrow \mathcal{L}f(x) + O(\epsilon)$$
 as $n \rightarrow \infty$
Computing transition rates

Let A, B $\subset \Omega$ be disjoint closed subsets of \mathcal{M} and N_{AB}(T) the number of transitions from A to B up to time T. Then using transition path theory (TPT) [3] the transition rate v_{AB} is given by

$$v_{AB} = \lim_{T \to \infty} T^{-1} N_{AB}(T) = \int_{\Omega} \|\nabla q\|_2^2 \mu(x) \operatorname{dvol}(x)$$

Here $\Omega = \mathcal{M} \setminus (A \cup B)$ and q is the committor function defined by the *committor BVP*:

(8) $\mathcal{L}q = 0$, $q|_{\partial A} = 0$, $q|_{\partial B} = 1$

TMD maps can be used as a meshless algorithm for numerically solving the committor problem since $L_{\epsilon,\mu}$ is a discrete approximation to \mathcal{L} :

$L_{\epsilon,\mu}q_{\mathsf{TMD}}(x) = 0 q_{\mathsf{TMD}} |_{\mathsf{A}\cup\mathsf{B}} = \mathbb{I}_{\mathsf{B}}$

References

- [1] Ralf Banisch, Zofia Trstanova, Andreas Bittracher, Stefan Klus, and Peter Koltai. "Diffusion maps tailored to arbitrary non-degenerate Itô processes". In: Applied and Computational Harmonic Analysis 48.1 (2020) pp. 242–265.
- [2] Ronald R Coifman and Stéphane Lafon. "Diffusion maps". In: Applied and *computational harmonic analysis* 21.1 (2006), pp. 5–30.
- [3] Eric Vanden-Eijnden and Weinan E. "Towards a theory of transition paths". In: Journal of statistical physics 123.3 (2006), pp. 503–523.

Sharp estimates for Target Measure Diffusion Maps show that convergence is faster when approximating the generator of an overdamped Langevin diffusion on its committor function with a quasi-uniform sampling density:

$$|4\beta^{-1}\mathsf{L}_{\epsilon,\mu}^{(n)}\mathsf{f}(\mathsf{x}) - \mathcal{L}\mathsf{f}(\mathsf{x})| = \frac{\epsilon}{(2\pi)^{d/4}} \sqrt{\frac{\log n}{\rho(\mathsf{x}_i)n\epsilon^{4+d/2}}} (2\|\nabla_{\mathcal{M}}\mathsf{f}(\mathsf{x}_i)\|\epsilon^{1/2})$$

The expressions for $\mathcal{B}_1, \mathcal{B}_2, \mathcal{B}_3$ are given by:

 $\mathcal{B}_2[f, \mu, \rho]$

 $\mathcal{B}_3[f,\mu,\rho]$

*B*₁[f,μ] :=

When less is more: Removing points to enhance spatial uniformity (δ -nets) improves the robustness of diffusion-map based transition rates for conformational changes in molecules!







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Let the i.i.d samples $\mathcal{X}(n)$, bandwidth ϵ , and target measure μ be used for constructing the TMD map generator $L_{\epsilon \mu}^{(n)}$. Then for $f \in C^2(\mathcal{M})$ and large enough n and small enough ϵ , with probability greater than $1-2n^{-3}$ we have:

$$\begin{aligned} \mathbf{J} &:= \frac{1}{16} \left(2\nabla \mathbf{f} \cdot \nabla \left(\mu^{1/2} \frac{\Delta \rho}{\rho} \right) + \left(\mu^{1/2} \frac{\Delta \rho}{\rho} \right) \mathbf{f} \right) \\ \mathbf{J} &:= \frac{1}{16} \left[\frac{\Delta (\mu^{1/2})}{\mu^{1/2}} - \left(\frac{\Delta \rho}{\rho} - \omega \right) \right] \left[\mathbf{f} \frac{\Delta (\mu^{1/2})}{\mu^{1/2}} - \frac{\Delta (\mu^{1/2} \mathbf{f})}{\mu^{1/2}} \right] \\ &= \frac{1}{4} \left[\mathcal{Q} \left(\mathbf{f} \mu^{1/2} \right) - \mathbf{f} \mathcal{Q} (\mu^{1/2}) \right] + \frac{\epsilon}{16} \left(2\nabla \mathbf{f} \cdot \nabla \left(\mu^{1/2} \omega \right) + (\mu^{1/2} \omega) \right) \right] \end{aligned}$$

Figure: (Left to right) The butane molecule $C_4 H_{10}$ can be effectively coarse—grained with the dihedral angle in the carbon backbone; $C_4 H_{10}$ is stable around $\theta \approx \pi$ and metastable around $\theta \approx \pi/3$, $5\pi/3$ and the intermediate angles can be sampled using metadynamics, with δ —net in \mathbb{R}^{12} further uniformizing the distribution; computing v_{AB} using TMD map shows that indeed this uniformization aids in improving the estimation of the transition rate and stabilizing it to choice of the bandwidth.





The key strategy to obtaining the prefactors is the expansion of integrals of the form

$$\mathcal{K}_{\epsilon}\rho(\mathbf{x}) := \int_{\mathcal{M}} \mathbf{k}_{\epsilon}(\mathbf{x}, \mathbf{y})\rho(\mathbf{y}) \, d\mathbf{y}$$

We use up to *fourth order* Taylor expansions of \mathcal{K}_{ϵ} in normal coordinates [2]:



 $\|s(u) - s(0)\|_{2}^{2} = \|u\|_{2}^{2} + \mathcal{C}_{x}^{1}(u), \mathcal{C}_{x}^{1}(u) = O(\|u\|_{2}^{4})$ $dy(u) = 1 + C_x^2(u), C_x^2(u) = O(||u||_2^2)$ Using the **discrete maximum principle** we prove an error estimate for solutions to Dirichlet BVP's:

Corollary. There exists a constant $C(\mathcal{M})$ such that with probability greater than $1 - 2n^{-2}$,

 $\|q_{\mathsf{TMD}}(\mathsf{x}_{\mathsf{i}}) - q(\mathsf{x}_{\mathsf{i}})\|_{\infty} \leq C \|4\beta^{-1}\mathsf{L}_{\varepsilon}^{(n)}q - \mathcal{L}q\|_{\infty}.$

Computing v_{AB} for C_4H_{10} with TMD map in 12-D,

 $A = \theta^{-1}(B_{0,2}(\pi)),$

 $\mathsf{B} = \theta^{-1}(\mathsf{B}_{0.1}(\pi/3) \sqcup \mathsf{B}_{0.1}(5\pi/3))$

8*	v_{AB}
0.0016	0.0109
0.0039	0.0112
N/A	0.0114
	E * 0.0016 0.0039 N/A